# **''Chimie Douce'' Synthesis and ab Initio Structure Determination** of (H<sub>3</sub>O)Yb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O: A Diamond Type Stacking of UOA<sub>[8]</sub>

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## **SYNTHESIS**

The fluoride  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  (SG:  $Fd\overline{3}m$ ,  $a=15.326(1)$  Å,  $Z=16$ ) was synthesized in the form of powder by a "chimie douce" route. The compound  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O (SG: *Fd* $\overline{3}m$ ,  $a=15.339(1)$  Å,  $Z=16$ ), a new member of the KYb<sub>3</sub>F<sub>10</sub> family, was obtained by an acid**–**base neutralization. The structural description is based upon the stacking of  $[Yb_6F_{32}]^{14-}$  "octahedral units of antiprism," called  $UOA_{[8]}$ , where the three-dimensional  $[Yb_6F_{20}]^{2-}$  network is simply a diamond type stacking of  $UOA<sub>181</sub>$  connected by vertices. The counterions are located in the middle of the chair hexagons. The zeolithic character of the water molecules is related to their mobility through the tunnels of this diamond structure. The simplifying description of this phase as a diamond type structure built up from  $UOA_{R}$  $(A_2[Yb_6F_{20}]$ ,  $2H_2O$ ,  $A=H_3O^+$ ,  $Z=8$ ) is the basis of our reexamination of the structural description of numerous complex structures built up of antiprisms.  $Q$  1997 Academic Press

## **INTRODUCTION**

Numerous fluorinated compounds involving lanthanide elements and alkaline metals, or alkaline earth metals, are known. This is not the case for fluorides in which 3d transition cations are associated with 4f rare earth ions. To our knowledge, only the MnLu $F_5$  compound (1) and the  $\text{Na}Ln\text{Cu}_2\text{F}_8$  family (2) are reported. For the f[orm](#page-9-0)er compound, no structura[l s](#page-9-0)tudy has been reported, but for the  $NaLnCu<sub>2</sub>F<sub>8</sub>$  series (*Ln* = Eu, Gd, Yb) the structure is claimed to be isotypic with the  $KBrF_4$  type. Our first aim was to prepare new fluorinated compounds with 4f and 3d elements, due to their potential magnetic properties. As all of our attempts in the solid state were unsuccessful, we decided to experiment with a ''chimie douce route'' via lanthanide and 3d oxalate precursors. During these trials, performed mainly with  $M = Mn$  and  $Ln = Yb$ , we found only the title compound. The present scope of this paper is to report the synthesis, the characterization, and the structure determination of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ .

A ''chimie douce'' process allowed for the synthesis of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ . In a first step, an amount of  $Yb_2$ was dissolved in a minimum amount of boiling nitric acid. An oxalic acid solution was then added to this solution so that a hydrated  $Yb_2(C_2O_4)$ <sub>3</sub>,  $xH_2O$  precipitate was obtained, then filtered out, washed with water, and air dried. In a second step, this salt was added to a heated solution of HF 49% in order to synthesize the fluoride. The most efficient temperature which yielded the best crystallinity of the resulting product  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  was empirically found to be 70*°*C. Our compound was obtained in the form of a fine white powder with no single crystal included.

#### SPACE GROUP DETERMINATION

X-ray powder patterns used for the structural determination were recorded on a Philips PW1380 goniometer equipped with a back graphite monochromator. Indexing of the first 25 reflections of the powder pattern was obtained utilizing the program TREOR (3), in a cubic system, with  $a = 15.322(1)$  Å. This solution [w](#page-9-0)as characterized by the following figures of merit:  $M(25) = 32$ ,  $F(25) = 35$ . Then, 72 reflections were indexed up to large angle with the cell parameter refined by the program ERACEL (4) remaining unchanged at 15.327(1) Å. The *hkl* observed re[fle](#page-9-0)ctions lead to the c.f.c. lattice. The highest symmetry space group *Fd*31*m* (no. 227, origin choice 2 at center  $\overline{3}m$ ) was deduced from the following conditions:  $hk0: h + k = 4n$  and  $h, k = 2n$ ;  $hhl : h + l = 2n; h00 : h = 4n.$ 

## CHARACTERIZATION

The fluorine analysis confirmed that the ionic network was made up completely of fluoride, as the mean experimental percentage,  $25.5(7)\%$  (2 measures), was in agreement with the 25.46% theoretical value for  $(H_3O)Yb_3F_{10}, H_2O$ . One can notice that if an anionic site had been occupied by  $OH^-$  instead of  $F^-$ , i.e., a 32e site (see crystallographic

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**FIG. 1.** DTA/TGA curves for  $(H_3O)Yb_3F_{10}, H_2O$ .

data), the fluorine percentage for  $(H_3O)Yb_3F_8(OH)_2, H_2O$ should have been much lower (20.48%).

Numerous data were collected from DTA (T.A. Instruments SDT 2960) and TGA (Setaram TGA 92 12) experiments. The coupled thermogravimetric and thermal analysis data collected at 400*°*C are shown in Fig. 1. It clearly exhibits a two-step loss of matter associated with three thermal phenomena. The first two peaks, both below 260*°*C, correspond to a loss of water without destruction of the crystalline building. The third exothermic peak between 275 and 330*°*C corresponds to the decomposition of the

starting compound to  $YbF_3$  with a loss of HF. These observations have been verified by X-ray thermodiffractometry analysis, shown in Fig. 2. The decomposition begins at 240 and ends at 360*°*C (the temperature difference between the thermodiffractometric and thermogravimetric techniques is approximately 20*°*C).

The first endothermic peak is correlated to a loss of weight which corresponds to zeolithic water. This phenomenon has been observed in a thermogravimetric study by which we heated our compound to a temperature of 130*°*C and then slowly cooled it to room temperature. It has been shown in a first analysis under argon that the loss of weight due to the heating was stabilized upon cooling. For the same experiment under air, we observed an increase of weight upon cooling, corresponding to a reabsorption of water (see Fig. 3).

A quan[titative](#page-2-0) analysis of our thermal studies is critical. The main objective is to determine the true amount of absorbed water due to time-dependent natural desorption. The total weight loss before decomposition reached 9.19% for a freshly synthesized sample (see Fig. 1). Two days later, the weight loss decreased to 8.02%. Finally, after a few weeks under atmospheric conditions, the observed loss was only 7.60%. The first two weight loss values indicated 1.5 to  $2 H<sub>2</sub>O$  per formula. The last value is close to the theoretical weight loss for  $(H_3O)Yb_3F_{10}, H_2O$  (7.51%) and corresponds to the stabilized water content which exhibits a zeolithic character. Indeed, as can be seen in Fig. 3, the weight increase (0.73 mg) when cooled under a[ir corr](#page-2-0)esponds to



**FIG. 2.** X-ray thermodiffraction of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  (vertical bars are related to YbF<sub>3</sub>).

<span id="page-2-0"></span>

**FIG. 3.** TGA experiment under air for  $(H_3O)Yb_3F_{10}, H_2O$ .

 $4.05 \times 10^{-5}$  mol of reabsorbed H<sub>2</sub>O. The mass of anhydrous fluoride,  $(H_3O)Yb_3F_{10}$ , at the end of the step was equal to 28.99 mg (30.21–1.22 mg), i.e.,  $3.98 \times 10^{-5}$  mol. The ratio of zeolithic water is thus close to one.

From a qualitative point of view, we can retain that three steps appear during the thermal treatment:

(1) zeolithic water desorption

$$
(H_3O)Yb_3F_{10}, H_2O \to (H_3O)Yb_3F_{10}+H_2O
$$

(2) elimination of  $H_2O$  associated to  $H_3O^+$ 

$$
(\mathrm{H_3O})\mathrm{Yb_3F_{10}} \rightarrow \text{``HYb_3F_{10}"} + \mathrm{H_2O}
$$

(3) decomposition

$$
HYb_3F_{10} \rightarrow HF + 3YbF_3.
$$

Steps (2) and (3) are not separable on the thermogravimetric curve, which is in agreement with the X-ray thermodiffractograms that show the appearance of  $YbF_3$  before the total disappearance of the starting compound. However, each step is distinguishable by TDA: (1) and (2) are endothermic phenomena, (3) is an exothermic decomposition.

DTA/TGA measurements show that the appearance of "HYb<sub>3</sub> $F_{10}$ " at the second step begins above 130<sup>°</sup>C. It is associated with the appearance of the low intensity 222 reflection at 120*°*C that is visible on the thermodiffractogram near  $20^{\circ}$  (2 $\theta$ ), and with a perceptible decrease of several intensities (for example: 111 reflection). Attempts to fit the 222 reflection from the " $HYb_3F_{10}$ " composition were unsuccessful.

An infrared spectroscopy study also confirmed the presence of H<sub>2</sub>O or  $H_3O^+$  molecules: a weak but broad band at  $1627$  cm<sup>-1</sup> and a very broad, intense band in the 3500-cm<sup>-1</sup> region were observed.



**FIG. 4.**  $\chi^{-1} = f(T)$  curve for  $(H_3O)Yb_3F_{10}, H_2O$ .

Finally, susceptibility measurements were performed in the temperature range 4.2*—*300 K using a Faraday method  $(H = 0.7$  T). The thermal variation of the reciprocal magnetic susceptibility (gram) is given in Fig. 4. A Curie*—*Weiss law is observed above 40 K. For the composition  $(H_3O)Yb_3F_{10}, H_2O$ , the calculated Curie constant and the paramagnetic Curie temperature were  $C_M = 8.26$ ,  $\theta p =$  $-42$  K, respectively. The effective moment per ytterbium ion, 4.69  $\mu_B$ , is then close to the theoretical 4.54  $\mu_B$  value for a Yb<sup>3+</sup> 5.09  $\mu_B$ , is then close to the theoretical 4.54  $\mu_B$  value for a 10<br>ion (*S* = 1/2, *L* = 3, *J* = *L* + *S* = 7/2,  $\mu_{\text{theo}} = g \sqrt{J(J+1)}$ with  $q = 8/7$ ). Below 40 K, the slight deviation from linearity is probably the consequence of the zero-field splitting of the Yb3*`* multiplet induced by the crystal field rather than the reflect of antiferromagnetic couplings.

#### AB INITIO X-RAY STRUCTURE DETERMINATION

A profile matching mode of the program FULLPROF (5) was used to calculate the intensities. The quality of t[he](#page-9-0) refinement was improved by suppressing the 111 strong reflection at  $10.01^{\circ}2\theta$ , which was highly asymmetric. The fit of the 422 and 822 reflections remained incorrect, the calculated intensities being inferior to the observed data. From this refinement ( $Rp = 11.9\%$ ,  $Rwp = 14.4\%$ ,  $\chi^2 = 30.2$ ,  $R_{\rm B} = 6.2\%$ ,  $R_{\rm F} = 3.7\%$ ), 196 intensities were converted into structure factors and used as input for the program SHELXS-86 (6). Scattering factors for  $Yb^{3+}$ ,  $F^-$ , and  $O^{2-}$ were taken fr[om](#page-9-0) the ''International Tables for Crystallography'' (7). The position of the heavy atom was located by using [th](#page-9-0)e option TREF, for 50 nonoverlapping reflections (at least  $0.06^{\circ}2\theta$  apart).

A refinement of the atomic coordinates and of the thermal motion of a Yb<sup>3+</sup> ion situated in a 48f site  $(\frac{3}{8} \frac{3}{8} 0.053)$ , <span id="page-3-0"></span>yielded to an agreement factor of  $R = 19.4\%$ . The difference-Fourier map showed several peaks with distances to heavy atoms characteristic of Yb*—*F bond length. The reliability factor of successive atomic position and thermal motion refinements of F<sub>1</sub> (96*e*: 0 0.87–0.87), F<sub>2</sub> (32*e*:  $x = 0.21$ , and F<sub>3</sub> (32*e*:  $x = 0.053$ ) decreased to 6.6%.

The deduced composition,  $Yb_{48}F_{160}$ , was not balanced as it corresponded to  $[Yb_3F_{10}]^-$ , with  $Z = 16$ . At this stage, the surrounding of the  $Yb^{3+}$  cation is well defined and corresponds to an almost regular antiprism  $YbF_8$ . The analysis of the difference-Fourier map allowed for the localization of an electronic density in the  $16d$  ( $\overline{3}m$ ) site. We assigned this density to an  $H_3O^+$  molecule (formally,  $O^{2-}$ , hydrogen atom being not visible). This counterion ensures the charge balance of the deduced  $(H_3O)Yb_3F_{10}$  formula. In the subsequent refinement, the thermal motion of the hydronium ion converged to  $4.1 \text{ Å}^2$  and the reliability factor decreased to 5.6%. Furthermore, an electronic density situated in a 48*f* site  $(\frac{3}{8}, \frac{3}{8}, 0.20)$  could be seen on the difference-Fourier map. This peak could be explained by additional water molecules. The occupation ratio of an  $O^{2-}$  ion situated in this site decreased to a third of its initial value and was then exactly fixed to  $\frac{1}{3}$ . At this condition, its thermal motion drops to about  $3 \text{ Å}^2$ , and the reliability factor decreased to 5.0%. The final composition was then  $(H_3O)Yb_3F_{10}, H_2O.$ 

At this point, we went back to the Rietveld refinement of the powder diffractogram. This calculation allowed for a better refinement of the position of the water molecules: it converged to  $\frac{3}{8}$   $\frac{3}{8}$  0.245, which generated H<sub>2</sub>O–F<sub>1</sub> distances

TABLE 1 Condition of Powder Data Recording (300 K) and Refinement for  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ 

| Angular range (°2 $\theta$ ):     | $5 - 40$                         |  |               |
|-----------------------------------|----------------------------------|--|---------------|
| Step (°2 $\theta$ ):              | 0.02                             |  |               |
| Time/step (sec):                  | 33                               |  |               |
| Receiving slits $(°)$ :           | 0.1                              |  |               |
| Refinement program:               | <b>FULLPROF</b>                  |  |               |
| Space group:                      | Fd3m(227)                        |  |               |
| Cell parameters (Å):              | $a = 15.326(1)$                  |  |               |
| Volume $(\AA^3)/\mathbb{Z}$ :     | 3599.9(7)/16                     |  |               |
| Zero point: $(^{\circ}2\theta)$ : | 0.0323                           |  |               |
| Background:                       | selected by hand                 |  |               |
| Number of reflections:            | 197                              |  |               |
| Peak shape:                       |                                  | Thompson–Cox–Hastings pseudo-Voigt                   |               |
| FWHM parameters:                  |                                  | $U_1 = 0.018(1)$ $V_1 = -0.024(2)$ $W_1 = 0.0101(7)$ |               |
| Asymmetry parameters:             | $0.173(3)$ $0.051(1)$            |  |               |
| Microstrain parameter:            | $Y = 0.101(1)$                   |  |               |
| Preferred orientation:            | $G1 = 0.940(3)$ (March function) |  |               |
| Number of parameters:             | 17                               |  |               |
| Conventional reliability          |                                  | $Rp = 14.2\%$ $Rwp = 16.9\%$ $\gamma^2 = 42.6$       |               |
| factors:                          |                                  | $R_{\rm B} = 7.56\%$                                 | $Rf = 6.11\%$ |
|                                   |                                  |  |               |

*Note*. Data recorded using a Philips PW 1380 diffractometer, Cu*K*a.

of 2.66 Å, instead of a shorter 2.2 Å distance. A 422 preferred orientation can be used to lower the  $I_{obs} - I_{cal}$  difference for the 422 and 822 reflections, but the problem concerning these two profiles remains unsolved. The conventional reliability values of the last refinement are  $Rp = 14.2\%$ ,  $R_{\text{WP}} = 16.9\%, \chi^2 = 42.6, R_{\text{B}} = 7.6\%, R_{\text{F}} = 6.1\%$ . The conditions of the powder data collection and refinement are reported on Table 1. The observed and calculated patterns



FIG. 5. Observed  $(\cdots)$  and calculated  $(\cdots)$  X-ray powder pattern of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ . The difference pattern is shown below at the same scale (vertical bars are related to the line's position).

TABLE 2 Atomic Parameters and Isotropic Beq Temperature Factors  $(A<sup>2</sup>)$  for  $(H<sub>3</sub>O)Yb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O$  (*Fd* $\overline{3}m$ , Origin: 16*c*)

| Atom             | Site    | X                           | ν             | z             | Beq     |
|------------------|---------|-----------------------------|---------------|---------------|---------|
| Yb               | 48 f    | $\frac{3}{8}$               | $\frac{3}{8}$ | 0.0531(1)     | 0.59(1) |
| $F_1$            | 96h     | $\theta$                    | 0.8749(4)     | $-0.8749(4)$  | 0.29    |
| F <sub>2</sub>   | 32e     | 0.2118(4)                   | 0.2118(4)     | 0.2118(4)     | $0.22*$ |
| $F_{3}$          | 32e     | 0.0538(4)                   | 0.0538(4)     | 0.0538(4)     | $0.22*$ |
| $H_3O^+$         | 16d     |                             | $\frac{1}{2}$ | $\frac{1}{2}$ | 3.7(7)  |
| H <sub>2</sub> O | $48f^a$ | $\frac{1}{2}$ $\frac{3}{8}$ | $\frac{3}{8}$ | 0.245(2)      | $3.7*$  |

*a* Occupation ratio  $\frac{1}{3}$ .

3 *\** denotes fixed factor.

are presented in Fig. 5. The atomic parameters and thermal motions are giv[en on T](#page-3-0)able 2.

## TEM ANALYSIS

An electron microscopy study confirmed the structure determination. An EDX analysis, performed with a JEOL-2010 TEM equipped with a KEVEX energy dispersive Xray spectrometer, revealed crystals with the mean composition  $Yb_3F_{10}O$ . No impurity was detected. Cell parameters were found to be  $a = b = c = 15.3(2)$  Å. The deduced reflection conditions:

— *hkl* : *h*, *k*, *l* same parity (*F* lattice)

 $\mu$  - 0*kl*:  $k + l = 4n$ , and  $k, l = 2n$  (*d* type plane)

—  $00l : l = 4n$ 

lead to the following possible space groups:  $Fd\overline{3}m$  (or  $Fd\overline{3}$ ) and  $Fd\overline{3}c$ . A careful examination of the intensities showed the existence of a multiple diffraction phenomenon: it was found that the intensity of the *hhl* with  $h + l = 2n$ , but  $h, l \neq 2n$  reflections was constant regardless of the tilt of the crystal. Thus, only the  $hhl : h + l = 2n$  condition was perfectly confirmed. The  $Fd\overline{3}c$  space group was then discarded. Finally, the deduced extinction symbol was  $Fd - (Fd3m)$  or  $Fd\overline{3}$  space groups), which is in perfect agreement with the X-ray diffraction study.

Figure 6a shows an electron diffraction pattern of a twinned crystal. Such patterns have often been observed. By comparison to a pattern of a nontwinned crystal (Fig. 6b) with the same zone axis, the first pattern can be explained by the twinning of two crystals having the same [101] zone axis with a  $\{111\}$  twinning plane, the twinning angle being 109*°*5. The [111]*\** row is common to the two crystals, as well as the [101]*\** row which is parallel to the incident beam.

HREM images of a complete focal series could not be taken because of a fast deterioration of the crystal under the beam. A simulation has been performed with the help of the program EMS of Stadelman (8) by the multislice method. The agreement between the o[bse](#page-9-0)rved and calculated image (Fig. 7) is very good.

#### STRUCTURAL DESCRIPTION

The coordination polyhedron of the  $Yb^{3+}$  ion is a square antiprism involving eight F<sup>-</sup> (see Fig. 8). The mean Yb-F distance,  $2.25(1)$  Å, is in a good ag[reemen](#page-5-0)t with the theoretical distance  $(2.28 \text{ Å})$  for eight coordinated Yb<sup>3+</sup> ions (Table 3). The maximum and minimum Yb*—*F distances  $(2.211(6)$   $\AA < d < 2.345(6)$  Å) are close to the observed extrema for  $\beta$  and  $\gamma$ -KYb<sub>3</sub>F<sub>10</sub> ( $\beta$ -KYb<sub>3</sub>F<sub>10</sub> (9): 2.108(5)  $\AA < d < 2.335 \AA$ ,  $\langle d_{\text{Yb-F}} \rangle = 2.23 \AA$  $\langle d_{\text{Yb-F}} \rangle = 2.23 \AA$  $\langle d_{\text{Yb-F}} \rangle = 2.23 \AA$ ;  $\gamma$ -KYb<sub>3</sub>F<sub>10</sub> (10): 2.199(5)  $\AA < d < 2.376 \AA$ ,  $\langle d_{\text{Yb-F}} \rangle = 2.29 \AA$ ). The [squ](#page-9-0)are faces of the antiprism are of different dimensions, as



**FIG. 6.** Electron diffraction pattern for  $(H_3O)Yb_3F_{10}, H_2O$ : (a) twinned crystal (b) nontwinned crystal.

<span id="page-5-0"></span>

FIG. 7. Observed and simulated HREM images along the [001] zone axis.

is usually observed for rare earths in such a coordination: the square face made by the four  $F_1$  anions exhibits  $F_1-F_1$ the square face made by the four  $F_1$  anions exhibits  $F_1-F_1$ <br>distances (2.710(6) Å) greater than the  $F_2-F_3$  distances (2.446(11) Å) observed for the square face made by the  $F_2$  and  $F_3$  anions. The Yb<sup>3+</sup> ion is correlatively slightly offcenter, being closer to the largest face: Yb-F<sub>1</sub> distances are shorter than  $Yb-F_2$  and  $Yb-F_3$  distances.

Each polyhedron shares the four  $F_2-F_3$  edges of the smallest square face with four neighboring antiprisms. As can be seen in Fig. 9, six antiprisms are therefore connected around a distorted cubic cavity for which angles range between 78.3(6)*°* and 100.6(5)*°* instead of 90*°*. This forms an  $[Yb_6F_{32}]^{14}$  structural unit, which resembles the geometry of an octahedron. For this reason, as well as for clarity, this unit will be further called *octahedral unit of antiprisms*, noted by  $UOA_{[8]}$ . In this abbreviation, the index [8] is used for



FIG. 8. Coordination polyhedron for Yb<sup>3+</sup>.

TABLE 3 Main Distances ( $\AA$ ) in  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ 

| $Yb^{3+}$ polyhedron: [8]  | $F-F$ distances in the square antiprisms  |
|--|---|
| $4 \times Yb-F_1$ : 2.211(6)<br>$2 \times Yb-F_3$ : 2.251(6)<br>$2 \times Yb-F$ , : 2.345(6)<br>$\langle d_{\text{Yb-F}} \rangle$ = 2.25(1) Å<br>$d_{\text{Shannon}} = 2.28 \text{ Å}$ | $2 \times F_1 - F_1 = 2.709(6)$<br>$2 \times F_1 - F_1 = 2.710(6)$<br>$4 \times F_1 - F_2 = 2.892(7)$<br>$4 \times F_1 - F_3 = 3.063(10)$<br>$4 \times F_2 - F_3 = 2.446(11)$ |
| Center of the hexagonal cavity<br>$6 \times (H_3O^+)$ -F <sub>1</sub> :2.709(7)  |   |
| Zeolithic water<br>$4 \times (H_2O) - F_1$ : 2.66(1)   |   |

indicating the number of anions constituting the central cavity (cubic in this case).

Each free vertex of an octahedral unit of antiprisms is shared with another  $UOA_{[8]}$ . Moreover, two  $UOA_{[8]}$  are linked by six vertices, noted, a, b, c, d, e, f in Fig. 9. Each  $UOA_{[8]}$  is then linked to four neighboring  $UOA_{[8]}$  tetrahedrally (see Fig. 10a). The formula of the resulting threedimension[al struct](#page-6-0)ure is then  $[Yb_6F_{20}]^2$ :  $Yb_6F_{32} \equiv$ <br> $Yd_6F_3 = Yd_6F_4$  $\text{Yb}_6\text{F}_8\text{F}'_{24} \rightarrow \text{Yb}_6\text{F}_8\text{F}'_{24/2} \equiv \text{Yb}_6\text{F}_{20}$ . As can be seen in Fig. 10b, in which each  $UOA_{[8]}$  is symbolized by a sphere, [the three](#page-6-0)-dimensional network of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  can be described as a diamond type stacking of  $UOA_{[8]}$  sharing vertices. The true formula is thus  $(H_3O)_2 Yb_6F_{20}$ ,  $2H_2O$ , with  $Z = 8$ .

Tunnels along [110] directions (see Fig. 11), chair hexagons and cavities (see Fig. 12) are of [course t](#page-6-0)he main features of this noncompa[ct diamo](#page-6-0)nd type stacking. The  $\rm H_3O^+$ 



**FIG. 9.**  $[Yb_6F_{32}]^{14-}$  structural unit: an  $UOA_{[8]}$  whose empty central cavity is a cube.

<span id="page-6-0"></span>

**FIG. 10.** (a) structure of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  viewed along  $[100]:[Yb_6F_{20}]^2$  network (b) the same view with spheres instead of  $UOA_{[8]}$ .

cations are located in the middle of the chair hexagons [16*d* position), at the center of an open face of pseudo hexagonal symmetry, leading to six  $(H_3O)$ – $F_1$  distances equal to  $2.709$  Å. As can be seen in Fig. 13, the water molecules are statistically distributed ins[ide the](#page-7-0) large cavities centered in the 8*b* position  $(\frac{3}{8}, \frac{3}{8})$ . These cavities are made of 28 F<sup>-</sup>, the mean ((8*b* site)–F) distance being 4.29 Å (24*x* 4.284(6) Å  $+4x$  4.331(6) Å). The free radius of the sphere corresponding to this cage is 2.98 Å, its free volume being  $111 \text{ Å}^3$ . Four open faces of pseudo-hexagonal symmetry delimit this cage, each face being common to two cages.

Four  $(H_2O)$ -F<sub>1</sub> distances, 2.66(1) Å, are observed.  $(H_3O)$ – $F_1$  and  $(H_2O)$ – $F_1$  are thus in the range 2.56–2.86 Å quoted by Simonov and Butvesky (11) and characteristic for O-H ··· F hydrogen bonding in [me](#page-9-0)tal fluoride hydrates. Bond length analysis is also presented in Table 4. The  $F_1$  anion exhibits the lowest valency, as it is li[nked by h](#page-7-0)ydrogen bonding together to  $H_3O^+$  and  $H_2O$ .



FIG. 11.  $(H_3O)Yb_3F_{10}, H_2O$ : structure viewed along [110] showing tunnels.



FIG. 12.  $(H_3O)Yb_3F_{10}, H_2O$ : structure viewed along [111].

<span id="page-7-0"></span>

FIG. 13. Cage built up from the connection of 10  $UOA_{181}$ . The open faces are occupied by  $H_3O^+$  and the cavity, statistically, by zeolithic water molecules.

## CATIONIC EXCHANGE AND SYNTHESIS OF *δ*-KYb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O

The zeolithic properties of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  have been clearly related to the structure of this compound which exhibits a three-dimensional network of tunnels. As the  $H_3O^+$  cations are also situated in these tunnels, we thought it possible to replace them by an alkaline cation of equivalent radius. Since the  $H_3O^+$  effective radius was 1.41 Å (2.709*—*1.30), we decided to make experiments with potassium ( $r = 1.38$  Å in a sixfold coordination).

A first test of cationic exchange by mixing our fluoride with a solution of KF was unsuccessful. In a second trial, we tried to obtain a potassium precipitate by adding the starting oxalate salt to a heated HF/KF (5 *M*) solution. As a result, we synthesized a badly crystallized sample of cubic  $\gamma$ -(or  $\alpha$ -)KYb<sub>3</sub>F<sub>10</sub> whose refined cell parameter was  $a =$ 11.301(5) Å (*Note: the KYb*<sub>3</sub> $F_{10}$  *family exhibits three phases*:  $\alpha - 769$ <sup>*°C*</sup> $\rightarrow \beta$  (*hexagonal*)  $- 952$ <sup>*°C*</sup> $\rightarrow \gamma$  (*cubic*). *The*  $\alpha$  *and*  $\gamma$  *forms have been reported to be isotypic* (10). *The*  $\alpha \rightarrow \beta$  transformation is not observable by DTA, bu[t on](#page-9-0)ly by *electric conductivity*). The thermal behavior of our compound, studied both by DTA and X-ray thermodiffraction

TABLE 4 Bond Length Analysis in  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ 

| Ions           |  | $\sum s$ | $S_{\text{th}}$ |
|----------------|--|----------|-----------------|
| Yb             | $4 \times 0.407 + 2 \times 0.291 + 2 \times 0.368$ | 2.95     |                 |
| $F_1$          | $2 \times 0.407$                                   | 0.81     |                 |
| F <sub>2</sub> | $3 \times 0.291$                                   | 0.87     |                 |
| $F_{3}$        | $3 \times 0.368$                                   | 1.10     |                 |

techniques, was different from the observed behavior for the  $KYb_3F_{10}$  phases synthesized in the solid state at high temperature. Indeed, above 450*°*C, the compound partially decomposed to  $\delta$ -KYb<sub>2</sub>F<sub>7</sub>, then above 750<sup>°</sup>C to YbOOH +  $\delta$ -K<sub>3</sub>YbF<sub>6</sub>, and finally to Yb<sub>2</sub>O<sub>3</sub> above 900<sup>°</sup>C.

At last, an acid*—*base neutralization was performed successfully at room temperature under stirring according to the reaction

$$
\begin{aligned} \text{(H}_3\text{O)}\text{Yb}_3\text{F}_{10}, \text{H}_2\text{O} + \text{KOH}(0.5 \text{ M})\\ \rightarrow \text{K}\text{Yb}_3\text{F}_{10}, \text{H}_2\text{O} + 2\text{H}_2\text{O}. \end{aligned}
$$

The resulting sample is a diamond type compound, K*`* replacing  $H_3O^+$ . In order to avoid confusion, this new phase has been called  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>,H<sub>2</sub>O.

The crystallographic analysis of this compound has been determined by powder X-ray diffraction. The cell parameter  $a = 15.339(1)$  Å, is very close to that of  $(H_3O)Yb_3F_{10}, H_2O$ . The conditions of data recording on a Siemens D500 diffractometer (Cu*K*a) and refinement are reported on Table 5. The structure refinement of  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>,H<sub>2</sub>O has been performed by analogy with  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  (the only visible difference between these two diffractograms corresponds to the presence of the low but nonnegligible 222 reflection). The 111 reflection was also discarded because of a high asymmetry. Calculated intensities of the 422 and 822 reflections were inferior to the observed ones, as in the first case, but the  $I_{\text{obs}} - I_{\text{cal}}$  difference was far smaller. The conventional values of the last refinement are  $Rp = 11.2\%$ ,  $Rwp = 12.3\%, \chi^2 = 13.0, R_B = 6.99\%, R_F = 4.96\%.$  The observed and calculated patterns are presented in Fig. 14. A refinement of the occupation ratio of the  $K^+$  [position](#page-8-0) proved that this site was fully occupied, as  $\tau$  converged to 98%. It is also noticeable that the  $R_p$  or  $R_p$  values are about 5% greater when K<sup>+</sup> is replaced by  $H_3O^+$  (formally  $O^{2-}$ ).

TABLE 5 Condition of Powder Data Recording (300 K) and Refinement for  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O

| Angular range (°2 $\theta$ ):     | $7 - 127$              |  |                   |
|-----------------------------------|------------------------|--|-------------------|
| Step (°2 $\theta$ ):              | 0.03                   |  |                   |
| Time/step (sec):                  | 13                     |  |                   |
| Receiving slits $(°)$ :           | 0.1                    |  |                   |
| Space group:                      | $Fd\overline{3}m(227)$ |  |                   |
| Cell parameters (Å):              | $a = 15.339(1)$        |  |                   |
| Volume $(\AA^3)/\mathbb{Z}$ :     | 3609.0(7)/16           |  |                   |
| Zero point: $(^{\circ}2\theta)$ : | 0.0082                 |  |                   |
| Number of reflections:            | 173                    |  |                   |
| Peak shape:                       |                        | Thompson–Cox–Hastings pseudo-Voigt                   |                   |
| FWHM parameters:                  |                        | $U_1 = 0.036(1)$ $V_1 = -0.029(2)$ $W_1 = 0.0273(6)$ |                   |
| Asymmetry parameters:             | 0.112(4)               | 0.037(1)   |                   |
| Number of parameters:             | 18                     |  |                   |
| Conventional reliability          | $Rp = 11.2\%$          | $Rwp = 12.3\%$                                       | $\gamma^2 = 13.0$ |
| factors:                          |                        | $R_{\rm B} = 6.99\%$                                 | $Rf = 4.96\%$     |

*Note*. Data recorded using a Siemens D500 diffractometer, Cu*K*a.

<span id="page-8-0"></span>

**FIG. 14.** Observed ( $\cdots$ ) and calculated (——) X-ray powder pattern of  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>, H<sub>2</sub>O. The difference pattern is shown below at the same scale.

Atomic parameters and thermal motions are given on Table 6. The interatomic distances (Table 7) are similar to those of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ .

The thermal behavior of  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>,H<sub>2</sub>O has been studied by TGA/DTA. The exothermic decomposition of this phase is observed between 429 and 455*°*C (peak maximum at 434*°*C). Above this temperature, the anhydrous diamond type  $\delta$ -KYb<sub>3</sub>F<sub>10</sub> becomes mainly the cubic  $\alpha$ - $\text{KYb}_3\text{F}_{10}$  form, with an additional amount of YbF<sub>3</sub>. By comparison with  $(H_3O)Yb_3F_{10}, H_2O$ , the decomposition occurs at a higher temperature  $(+143^{\circ}C)$ . Similarly, the loss of zeolithic water between 150 and 250*°*C is also observed at higher temperature (this phenomenon was achieved at 130<sup>°</sup>C for the zeolithic water of  $(H_3O)Yb_3F_{10}$ ,  $H_2O$ ). This

difference may be due to the presence of the  $K^+$  ions in the center of the windows of the cages, which probably produce a more difficult desorption of the water molecules through the tunnels.

It is noticeable that the infrared spectra exhibit the same bands (1630 and 3400–3600 cm<sup>-1</sup>) as in the case of  $(H_3O)Yb_3F_{10}, H_2O$ . However the observed bands are less broad, in agreement with the absence of the vibration modes of  $H_3O^+$ .

### **CONCLUSION**

The original phase  $(H_3O)Yb_3F_{10}$ ,  $H_2O$  has been synthesized at low temperature by a ''chimie douce'' route. The



<span id="page-9-0"></span>synthesis of isotypic compounds of the lanthanide series is now in progress. A neutralization also allowed for the synthesis of  $\delta$ -KYb<sub>3</sub>F<sub>10</sub>,H<sub>2</sub>O which is a new compound of the  $KYb_3F_{10}$  family.

These two structures have been described in a simple fashion based on the stacking of groups of six antiprisms, connected by edges around a cubic cavity, and defined as an  $Yb_6F_{32}$  UOA<sub>[8]</sub>. The  $Yb_6F_{20}$  three-dimensional network corresponds to a diamond stacking of  $UOA_{[8]}$ connected by vertices. The counterions  $(H_3O^+, K^+)$  are located at the center of the faces delimiting the cages of the structure. The zeolithic character of the water molecules is related to their mobility through the tunnels of this diamond structure.

The greatest interest of this work is to provide a very simplifying structural description based on the modes of connection between polyhedra instead of a complex one. In fact, the scale of the basic unit of description has grown from an atom or a polyhedron to a group of polyhedra. From this observation, the members of the  $KYb_3F_{10}$  family and several structures based on antiprism have also been reconsidered. The results of this study are reported in the following publication.

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